

General announcements

Electric flux

We can describe the piece of paper (oriented parallel to the screen again) with an **area vector**

- An **area vector** points perpendicular to the surface whose area you're representing (a normal vector!)

So let's say the area vector for the piece of paper is pointing away from the screen (in the same direction as the electric field).

Electric flux (Φ_E) is produced by the **electric field component that is parallel to the area (normal) vector for a surface.**

- *Electric flux is* a way to measure how much of the electric field (e.g. the number of field lines) is passing through a surface of area A .

Mathematically:

$$\Phi_E = \vec{E} \cdot \vec{A} = EA \cos \theta$$

Units: $\text{N} \cdot \text{m}^2 / \text{C}$

The angle between E and A

What if the surface isn't regular?

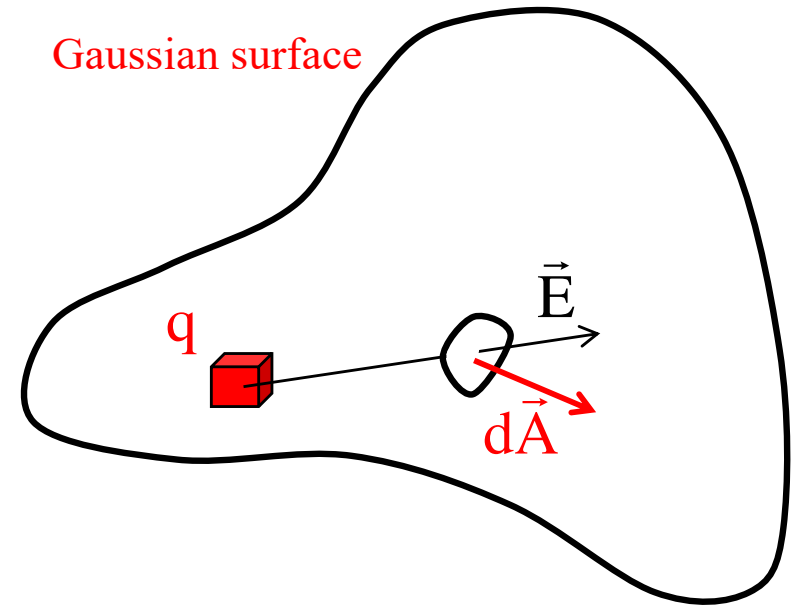
If you have an irregular surface, then simply dotting \vec{E} and \vec{A} won't work. This requires calculus!

We won't be evaluating these integrals, but it's worth seeing what this would be like in a non-ideal, real-world problem

Imagine a charge, q , inside an irregular surface as shown. At some differential area on the surface, $d\vec{A}$, the electric field is as shown.

Summing up all the $d\vec{A}$ patches, and the \vec{E} at each point (which is NOT uniform), would give us the net flux:

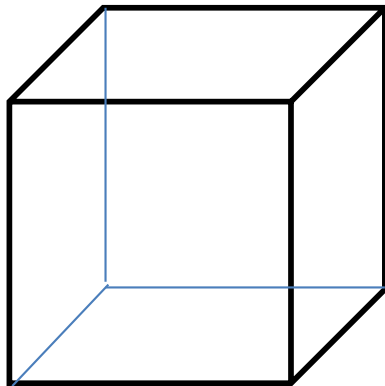
$$\int \vec{E} \cdot d\vec{A}$$



Flux through closed surfaces

Back to the cube: Let's say there's a charge Q to the left of the cube as shown. What's the net electric flux through the cube?

Q

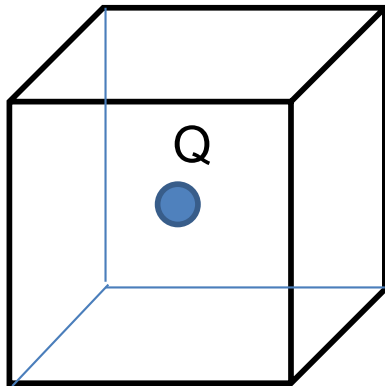


The field lines from Q will enter the cube on the left...but then exit on the right. Every line that enters the cube (which is negative flux) will also exit the cube (which is positive flux).

This means the net flux through the cube will be zero!

Flux through closed surfaces

What if the charge is now inside the cube?



The field lines from Q will now exit the cube through each face...and there will be a net positive electric flux through the cube.

Carl Gauss noticed this too, and realized that **as long as there is charge enclosed within the surface, there will be a net electric flux through that surface.**

Gauss's Law

Gauss realized the net electric flux through a closed surface is proportional to the net charge enclosed in that surface...but he needed a proportionality constant.

Imagine a point charge inside a sphere.

Using the idea of flux, and assuming E is perpendicular to A:

$$\Phi_E = \vec{E} \cdot \vec{A} = EA = k_e \frac{Q}{r^2} (4\pi r^2) = 4\pi k_e Q$$

Often, instead of using k_e , we use a constant called the **permittivity of free space**, ϵ_0 , which is equal to $\frac{1}{4\pi k_e} = 8.85 \times 10^{-12} \text{ C}^2 / (\text{Nm}^2)$

This makes our equation above into:

$$\Phi_E = \frac{Q_{\text{inside}}}{\epsilon_0}$$

This is Gauss's Law!

Gauss's Law

This form of Gauss's Law is only useful when you have a symmetric surface about the point charge producing the electric field.

This does not have to be a real surface! We are creating a mathematical, imaginary surface to aid in our calculations. These are called "Gaussian surfaces."

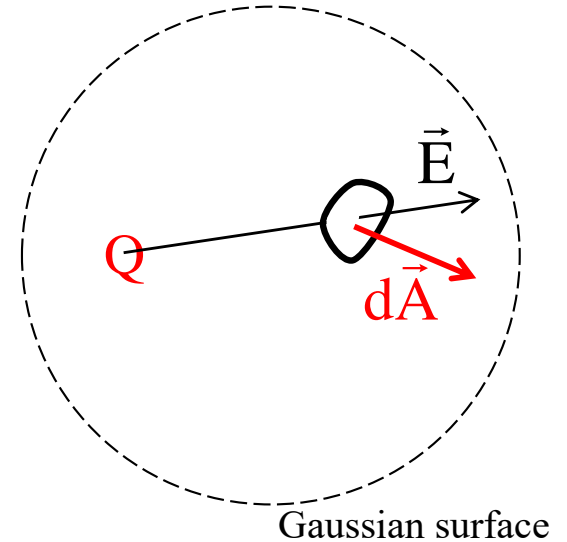
As long as we choose a surface (like a sphere) so that the electric field is constant everywhere on it, we can use this equation to compute the electric field.

In the next few slides, I present an example to show you why and how this is so powerful. Thanks, Mr Fletcher, for this set up!

Back to Gauss's Law

Gauss's Law is ALWAYS TRUE, no matter how the geometry shakes out, but it is pretty useless unless you can exploit symmetry in a problem.

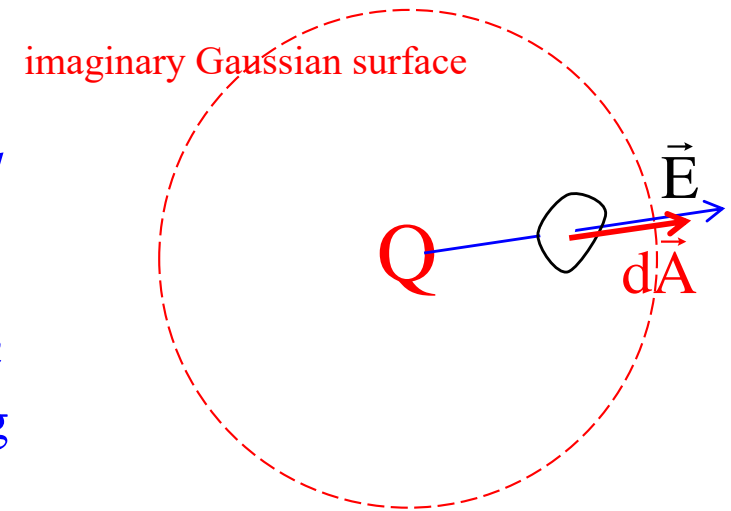
Example 1: Use *Gauss's Law* on the spherical surface of radius R and charge Q as shown.



What the integral $\int_S \vec{E} \cdot d\vec{A}$ is apparently asking us to do is to **define** an **arbitrary, differential surface area dA** (remember, dA has a magnitude equal to the area of the enclosed surface and is directed perpendicularly outward from the surface), **evaluate** both the **direction** and **magnitude** of the **electric field at that surface**, **dot the \vec{E} and $d\vec{A}$ into one another**, then do that for all the differential surfaces over the entire structure and **sum them by integrating**.

Gauss was right. That flux **WOULD** equal Q/ϵ_0 . But because **no two points on the surface are the same distance from the charge Q** , and **no two dot products are going to be the same (angles different)**, doing that integral would be a **NIGHTMARE**. In short, this is an impossible problem to do!!!!!!

Example 2: This is *Example 1* done more reasonable: **Derive an expression** for the *electric field function* for a point charge Q .

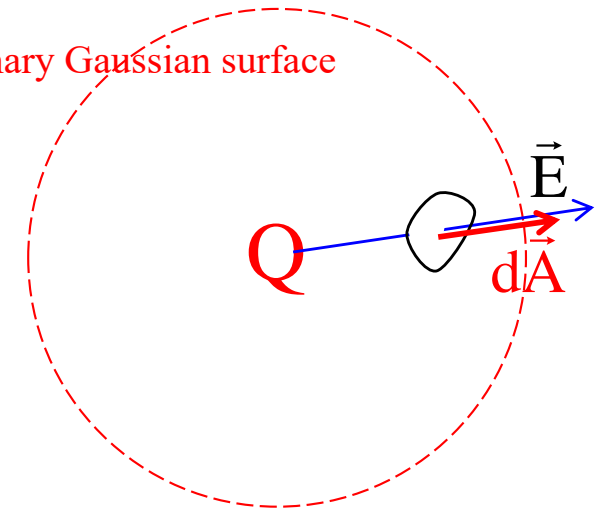


Important observation: There is *no given Gaussian surface* to begin with in this problem, just a hanging charge. We need to **create an imaginary Gaussian surface around the charge**, one that *exploits the symmetry* of the **charge's electric field**. That is, we need to **create a surface** such that *every point* on the surface is *equidistant from the charge*.

With the imaginary Gaussian surface centered on the charge, ANY differential area vector $d\vec{A}$ will be radially outward, which is to say, *in the direction of E* , and the **angle between the electric field vector and the differential area vector** will be **zero** (so the **cosine** in the *dot product* will equal **1**). With that, we can draw (see sketch), then write:

$$\int_S \vec{E} \cdot d\vec{A} = \frac{q_{\text{enclosed}}}{\epsilon_0}$$
$$\Rightarrow \int_S |\vec{E}| |d\vec{A}| \cancel{\cos 0^0}^1 = \frac{Q}{\epsilon_0}$$

imaginary Gaussian surface



Herein lies the beauty of the method. Because every point on the surface is equidistant from the charge, the evaluation of the E at every differential surface dA WILL BE THE SAME, which is to say, IS A CONSTANT VALUE, and because it is a constant, we can pull it out of the integral. (Note that we couldn't do that with the original Example 1 because each point was a different distance from Q .) With that, we can write:

$$\int_s |\vec{E}| |d\vec{A}| = \frac{Q}{\epsilon_0}$$
$$\Rightarrow |\vec{E}| \int_s |d\vec{A}| = \frac{Q}{\epsilon_0}$$

That makes life wonderful, as now the only thing inside the integral is the differential surface area dA , and summing that over the surface simply yields the total surface area of the sphere ($4\pi R^2$) . . . So we can further write

$$|\vec{E}| \int_s |d\vec{A}| = \frac{Q}{\epsilon_0}$$
$$\Rightarrow |\vec{E}| (4\pi R^2) = \frac{Q}{\epsilon_0}$$
$$\Rightarrow |\vec{E}| = \frac{Q}{4\pi\epsilon_0 R^2}$$

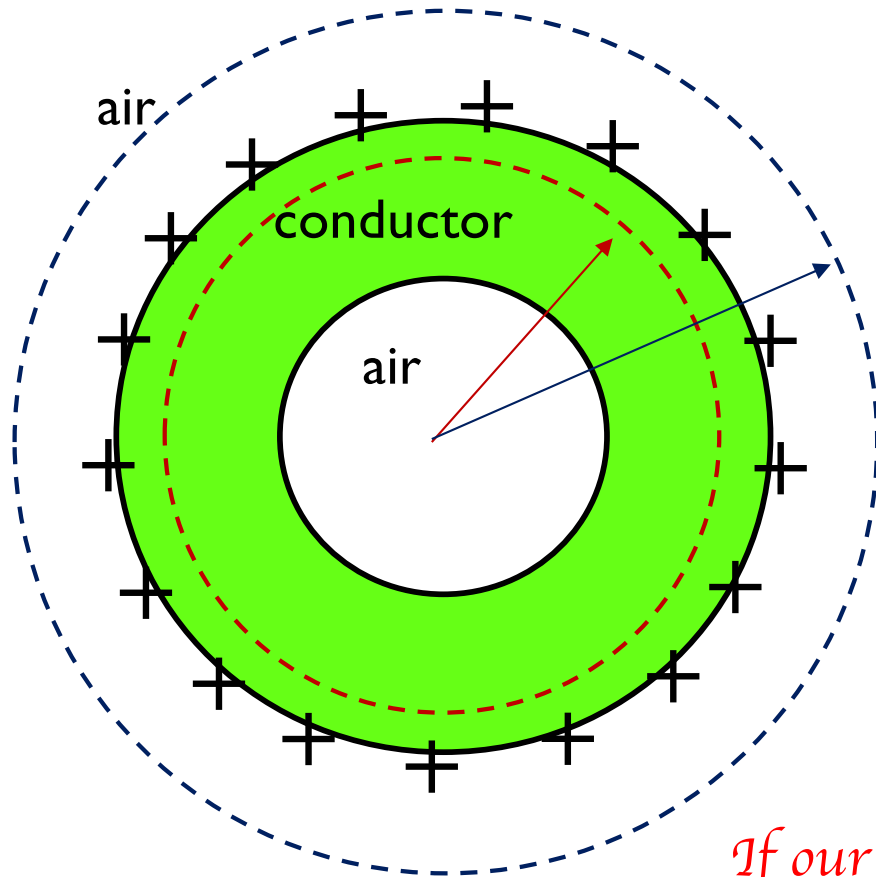
Look familiar? It should. It's the same as the electric field function we derived for a point charge using Coulomb's Law!

So why is this useful?

Gauss's Law helps us understand what's going on inside conductors and insulators, and supports what we said about electrostatic equilibrium before. To recap:

- *Remember that* the field inside a conductor itself is 0!
- *Charges on the surface* of a conductor can produce electric fields outside the conductor (we've seen this).
- *Placing a charge inside* the conductor can also produce an electric field outside the conductor!
- *To wit...*

What does the electric field look like here?



The conductor has a net positive charge Q , distributed on its outer surface. From before, we know that the electric field inside the conductor should be zero. Does Gauss's Law support this?

If our Gaussian surface is inside the conductor, it encloses no net charge.

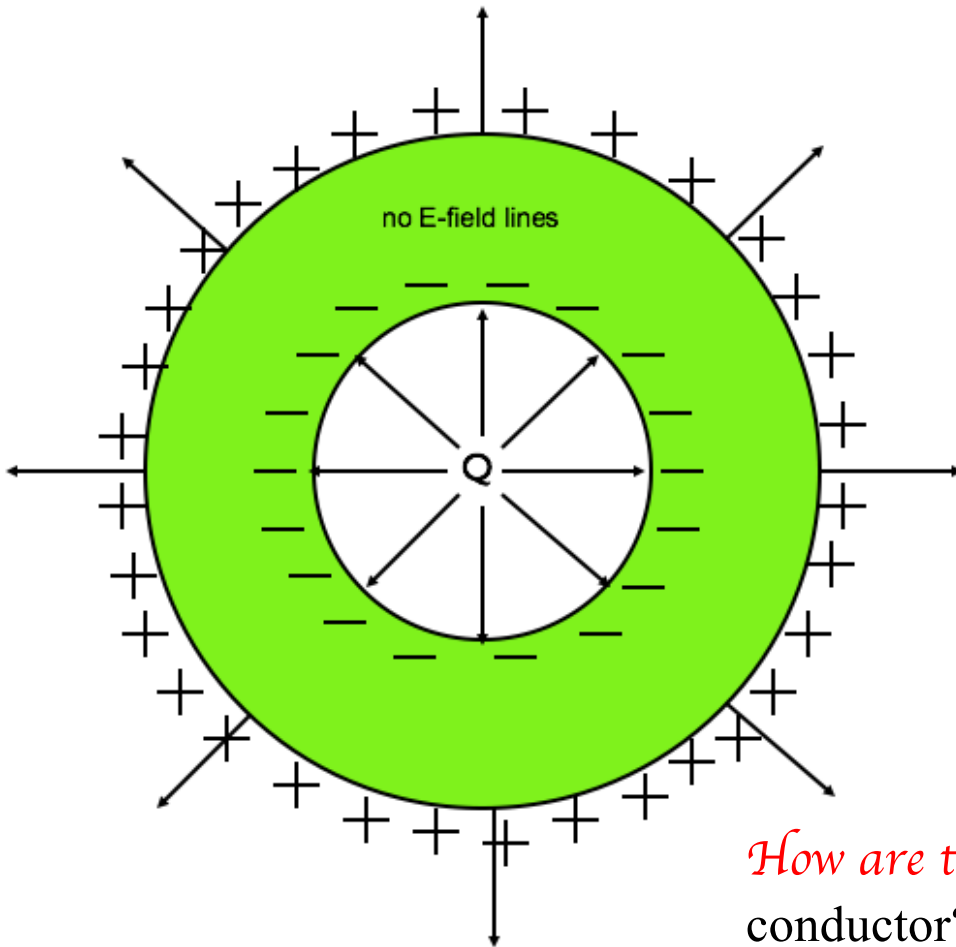
Because $\Phi_E = \frac{Q_{\text{inside}}}{\epsilon_0} = EA$, and $Q_{\text{inside}} = 0$, there is no electric field inside the conductor.

If our Gaussian surface is outside the conductor, it encloses charge Q . Thus, the electric field

$$\text{outside the conductor is } E = \frac{Q}{(4\pi r^2)\epsilon_0}$$

Another example

What do the electric field lines look like for a conducting sphere if there is a charge Q at its center?



This induces a $-Q$ charge on the inside surface of the conductor, and a $+Q$ charge on the outside.

Thus we end up with field lines as shown.

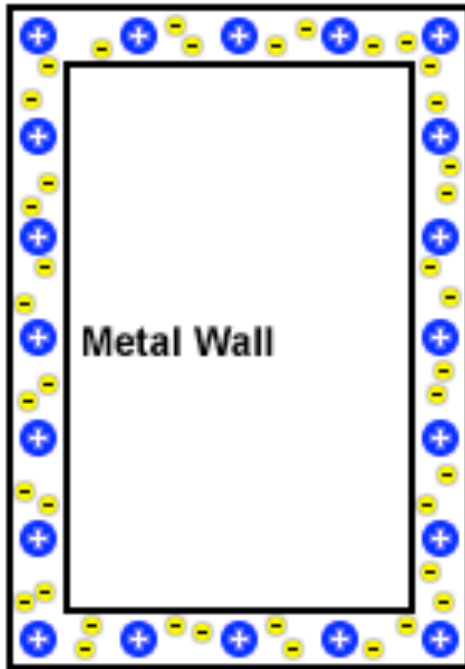
Why are there no field lines in the conductor? If our Gaussian surface is inside the conductor, it encloses no net charge!

How are there field lines outside the conductor?

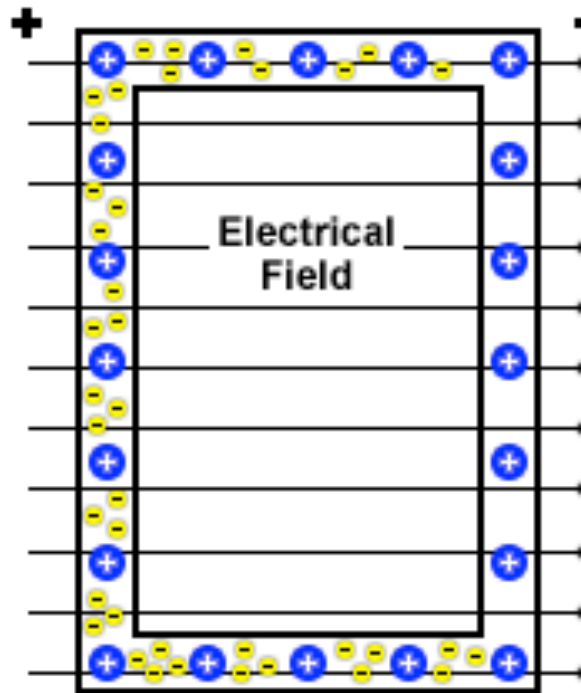
A Gaussian surface outside the conductor encloses a net charge of $+Q$!

Electrical shielding

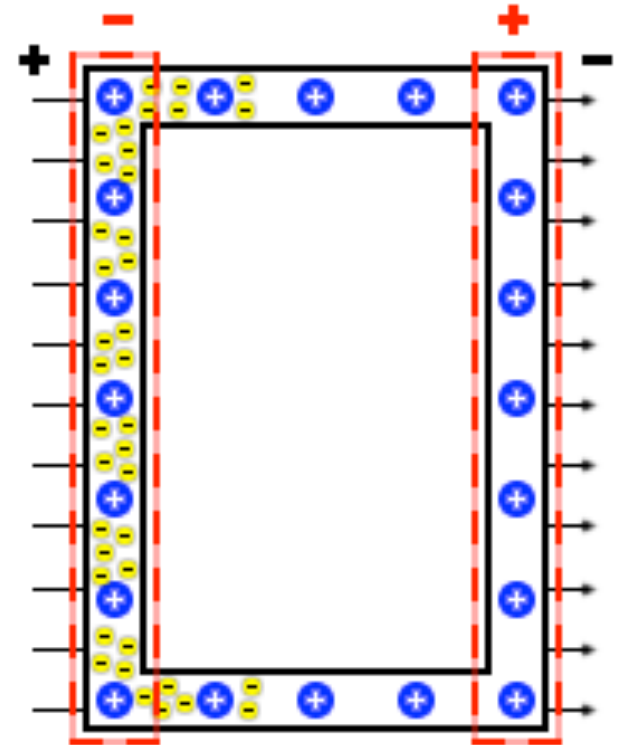
Faraday Cage



Faraday Cage in the absence of an electrical field.



The charged particles in the wall of the Faraday cage respond to an applied electrical field.



Electrical fields generated inside the wall cancel out the applied field, neutralizing the interior of the cage.

Faraday Cage

